Improving Robustness in Multiscale Methods

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Outline

- Modern Scientific Computing
- Multigrid Methods
- Upscaling and Homogenization
- Adaptive Multigrid Methods
- Future Work

Why Compute?

Interested in modeling physical processes

- Diffusion (Heat, Energy, Chemical)
- Fluid Flow
- Particle Transport
- Elastic Materials
- Can describe these processes through differential equations (both ODEs and PDEs)
- Cannot write down closed form solutions
- Need to find (approximate) solutions in other ways

Scientific Computation

- Significant interest in simulating complex physical systems with features, and hence solutions, that vary on multiple scales
- Accuracy constraints lead to discretizations with tens of millions, or even billions, of degrees of freedom (DOFs)
 - 3D Tsunami Model: 200 million cells
 - Transport: 500 million to 1 billion DOFs

Properties of Discretizations

- We consider discretizations of the underlying continuum models (differential equations) via finite elements (or finite differences)
- The matrices from these discretizations tend to be sparse and ill-conditioned
- The matrices inherit properties of the continuum model (e.g. symmetry, definiteness)

Classical Methods do not Suffice



Porous Media Flow

Saturated flow in a reservoir can be modeled via Darcy's Law:

$$\mathbf{u}(\mathbf{x}) = -\mathcal{K}(\mathbf{x})\nabla p(\mathbf{x})$$
$$\nabla \cdot \mathbf{u}(\mathbf{x}) = Q(\mathbf{x})$$

- Simulation domain: $\sim 10^3$ meters in each dimension
- Material properties ($\mathcal{K}(\mathbf{x})$) vary on millimeter scales

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- Simulation domain: $\sim 10^3$ meters in each dimension
- Material properties ($\mathcal{K}(\mathbf{x})$) vary on millimeter scales
- For 3-D, fully-resolved flow calculation, need $\sim 10^{18}$ DOFs
- Such a simulation is at the limits of the capability of modern supercomputers (the fastest of which performs 3.5×10^{13} floating point operations per second)

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 - Fully-resolved Porous Media Flow: 10^{18} DOFs
- Without optimal methods, solving such problems can be prohibitively expensive

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- Bottom Line:

Accuracy per computational cost

Stationary Iterative Methods

- The Jacobi and Gauss-Seidel iterations do converge for FE discretizations of elliptic operators, but require $O(N^2)$ operations for 2-D problems and $O(N^{\frac{5}{3}})$ operations for 3-D
- These methods do, however, resolve some components much faster than others
- For the Laplacian operator, it is the geometrically smoothest components of the solution that are the slowest to be resolved
- For this reason, Jacobi and Gauss-Seidel are often called smoothers they smooth the error in the approximation

Smoother Performance





Multigrid Methods achieve optimality through complementarity

Multigrid Methods achieve optimality through complementarity

- **Multigrid Components**
 - Relaxation

Relax • $A^{(1)}x^{(1)}=b^{(1)}$

Use a smoothing process (such as Gauss-Seidel) to eliminate oscillatory errors

Remaining error satisfies $Ae = r \equiv b - Ax$

Multigrid Methods achieve optimality through complementarity



Transfer residual to coarse grid

Multigrid Methods achieve optimality through complementarity



Use coarse-grid correction to eliminate smooth errors

To solve for error on coarse grid, use residual equation

$$A^{(2)}e^{(2)} = r^{(2)}$$

Multigrid Methods achieve optimality through complementarity

Multigrid Components

- Relaxation
- Restriction
- Coarse-Grid Correction
- Interpolation



Transfer correction to fine grid

Multigrid Methods achieve optimality through complementarity

Multigrid Components

- Relaxation
- Restriction
- Coarse-Grid Correction
- Interpolation
- Relaxation



Relax once again to remove oscillatory error introduced in coarse-grid correction

Multigrid Methods achieve optimality through complementarity

Multigrid Components

- Relaxation
- Restriction
- Coarse-Grid Correction
- Interpolation
- Relaxation



Obtain optimal efficiency through recursion

Geometric Multigrid

- When Ax = b comes from a geometrically regular discretization of a DE, that information can be used in the coarse-grid problems
- Coarse grids chosen by removing points from the fine grid in a geometrically regular fashion
- Restriction (R) and interpolation (P) operators computed using geometric locations
- Coarse-grid operators determined by rediscretization on the reduced space

BoxMG

- The Black Box Multigrid Algorithm (BoxMG) was developed by Dendy for discontinuous coefficient operators, such as -∇ · K(x)∇
- Coarsening is geometric
- Interpolation is chosen to approximately preserve continuity of normal flux
- Variational formulation

Variational Multigrid

- Multigrid with $R = P^T$ and $A_c = RAP$ is called a variational formulation
- Terminology comes from minimization form of Ax = b:

$$F(v) = \frac{1}{2} \langle Av, v \rangle - \langle b, v \rangle$$
$$x = \arg\min_{v \in \mathcal{H}} F(v)$$

Given an approximation, v, to the solution on the fine level, it can be shown that the optimal coarse grid correction, Pw, solves

$$(P^T A P)w = P^T (b - Av)$$

Coarse-Scale Models

- Quality of coarse-grid corrections depends on representation of fine-scale details in coarse-scale operator
- Rediscretization requires explicit averaging of fine-scale structures
- Variational coarsening allows multiscale information to be encoded in the coarse-grid operators
- Coarse-scale operator needs to reflect information about the fine-scale operator relating to the low-energy modes for which a coarse-scale correction is being computed.
- Practical interest is in coarse-scale properties, such as total net flux

The Need for Coarse-Scale Models

- Fine-scale discretization requires too many DOFs
- Coarse-scale properties are influenced by fine-scale variations
- Need physically-meaningful solutions
- Cannot directly (explicitly) capture effect of fine-scale variations in a coarse-scale discretization
- Goal: derive effective, coarse-scale models

Interpretation of Multigrid CGOs

Fine-scale, finite-element discretization of porous-media problem:

$$A_{ij} = \mathbf{e}_j^T A \mathbf{e}_i = \int_{\Omega} \langle \mathcal{K}(\mathbf{x}) \nabla \phi_i, \nabla \phi_j \rangle d\Omega$$

Variational coarsening gives finite-element discretizations on coarse grids:

$$(A_c)_{ij} = (P^T A P)_{ij} = (P \hat{\mathbf{e}}_j)^T A (P \hat{\mathbf{e}}_i)$$
$$= \int_{\Omega} \left\langle \mathcal{K}(\mathbf{x}) \nabla \left(\sum_l p_{li} \phi_l \right), \nabla \left(\sum_k p_{kj} \phi_k \right) \right\rangle d\Omega$$
$$= \int_{\Omega} \left\langle \mathcal{K}(\mathbf{x}) \nabla \hat{\phi}_i, \nabla \hat{\phi}_j \right\rangle d\Omega$$

Coarse-grid basis functions are linear combinations of fine-grid basis functions (weighted by the interpolation operators)

Periodic permeability field, $\mathcal{K}(\mathbf{x})$, with jump of 10^3



Bilinear basis function on $4\times 4~{\rm grid}$



 8×8 grid multiscale basis function



 16×16 grid multiscale basis function



 32×32 grid multiscale basis function



 64×64 grid multiscale basis function



Geostatistical permeability field, $\mathcal{K}(\mathbf{x})$, with range of $[10^{-2}, 10^2]$ (Black pixels correspond to $\mathcal{K} = 10^{-2}$)



Bilinear basis function on 4×4 grid







 16×16 grid multiscale basis function



 32×32 grid multiscale basis function



 64×64 grid multiscale basis function



Multiscale Goals

Capture macro-scale properties of fine-scale operator

- Modeling flow, so compute net flux through domain of interest
- Boundary conditions: p = 1 on left, p = 0 on right, no-flow (Neumann) on top and bottom
- Integrate $(\mathcal{K} \nabla p) \cdot \mathbf{n}$ along outflow
- Linear model is a simplification of unsaturated/multi-phase flow where $\mathcal{K} = \mathcal{K}(p, \mathbf{x})$
 - Local pressure fluctuations important
 - Seek to match local maxima and minima of p

Periodic Theory Approach

- Two-scale asymptotic analysis to compute homogenized permeability
- Upscaled equation, $-\nabla \cdot \hat{\mathcal{K}} \nabla p_0(\mathbf{x}) = Q(\mathbf{x})$, arises from solvability conditions with periodic BCs
- Upscaled permeability, $\hat{\mathcal{K}}$, given in terms of particular solutions of fine-scale problem over averaging subdomain
- Coarse-scale model created by rediscretization on that scale
- We've shown equivalence between this approach and that of Durlofsky

Flux Calculations

- Multigrid upscaling procedure
 - Discretize on fine scale
 - Use BoxMG to coarsen to given coarse scale
 - Solve coarse-scale problem
 - Interpolate solution to fine scale
 - Compute Outflow Flux
- Periodic theory approach
 - Discretize on fine scale
 - Solve two fine-scale problems per coarse-scale element to get $\hat{\mathcal{K}}$
 - Rediscretize on coarse scale
 - Solve coarse-scale problem
 - Compute flux with coarse-scale permeability and pressure

Flux Calculations

Computed Flux for 512×512 element discretization is 2.229

Coarse Grid	Multilevel Upscaling		Periodic Theory	
	Flux % Error		Flux	% Error
32×32	2.430	9.0%	2.319	4.0%
16×16	2.558	14.8%	3.482	56.2%
8×8	2.599	16.6%	4.923	120.8%
4×4	2.493	11.8%	3.124	40.1%

Fine-Scale Structure

Accurate reconstruction of fine-scale structure needed to address nonlinearities in unsaturated and multi-phase flows



Accuracy of Coarse-Scale Models

- Performance of multilevel upscaling technique comes from accuracy of coarse-scale model
- In variational multigrid setting, good coarse-scale models come from good choices in interpolation
- Good interpolation must complement relaxation

Algebraically-Smooth Error

- Multigrid methods reduce error through
 - Relaxation (Jacobi, Gauss-Seidel)
 - Coarse-grid correction (variational)
- Error which is not efficiently reduced by relaxation is called *algebraically* smooth and must be reduced by coarse-grid correction
- Pointwise relaxation implies that algebraically-smooth error, e, satisfies $Ae \approx 0$, relative to e
- If the origins of the matrix are known, so is character of algebraically-smooth error

Algebraic Multigrid

- Assume no knowledge of grid geometry
- Interpolation and coarse grids chosen based only on the entries of the matrix
- Primary goal is to interpolate suitable corrections from the coarse grids
- Assume algebraically-smooth error is locally constant
- Equivalently, assume global near null space is the constant vector

Adaptive Multigrid

If we don't know what algebraically-smooth error looks like, can we still develop an effective multigrid method?

Adaptive Multigrid

- If we don't know what algebraically-smooth error looks like, can we still develop an effective multigrid method? Yes!
- Use relaxation on Ax = 0 to expose algebraic smoothness
- Fine-grid relaxation quickly exposes local character of algebraic smoothness
- Use this representation to determine interpolation

Adapting Interpolation

- AMG and BoxMG choose interpolation by a fixed formula
- Now need interpolation that depends on both the matrix, A, and the prototypical algebraically-smooth error, v
- Algebraic smoothness still means that $Ae \approx 0$, or

$$a_{ii}e_i \approx -\sum_{k \in C_i} a_{ik}e_k - \sum_{j \in F_i} a_{ij}e_j$$

Approximate e_j by values in $C_i \cap C_j$, weighted by a_{jk} and v_k

Scaling Property

- If we scale $A \rightarrow DAD$ for diagonal matrix D, adaptive AMG performance need not suffer
- If $A\mathbf{v} = \mathbf{0}$, then $DAD(D^{-1}\mathbf{v}) = 0$
- Replacing v with $D^{-1}v$, the coarse-grid matrices for DAD are diagonally-rescaled versions of those for A
- Pointwise relaxation is also invariant to such scaling
- If we could generate D⁻¹v as easily v, overall performance wouldn't degrade

Test Problems

 $-\nabla \cdot \mathcal{K}(\mathbf{x}) \nabla p(\mathbf{x}) = 0 \text{ on } [0,1]^2$

Problem 1:

• $\mathcal{K}(\mathbf{x}) = 1$ (Laplace), Full Dirichlet BCs

Problem 2:

• $\mathcal{K}(\mathbf{x}) = 10^{-8}$ on 20% of elements chosen randomly,

 $\mathcal{K}(\mathbf{x}) = 1$ otherwise

Dirichlet BCs on left and right, Neumann on top and bottom

- Setup Phase: Single V-cycle, # pre-relaxations chosen for optimal performance
- Solution Phase: V(1,1) cycles until residual reduced by 10¹⁰ or 200 iterations
- Geometric choice of coarse grids

Numerical Results - Solution Time

	Problem 1		Problem 2		
h	Standard AMG	Adapted AMG	Standard AMG	Adapted AMG	
1/64	0.04s	0.04s	0.05s	0.03s	
1/128	0.22s	0.25s	0.28s	0.31s	
1/256	0.91s	0.89s	1.04s	1.09s	
1/512	3.32s	3.52s	4.40s	4.84s	
1/1024	13.13s	14.70s	17.64s	22.06s	

Total time to setup MG method and reduce residual by 10^{10}

Numerical Results - Convergence Factors

Asymptotic convergence	e factors of	f resulting \	V(1,1) cycles
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	Problem 1		Problem 2	
h	Standard AMG	Adapted AMG	Standard AMG	Adapted AMG
1/64	0.104	0.067	0.209	0.194
1/128	0.115	0.073	0.212	0.202
1/256	0.124	0.079	0.233	0.243
1/512	0.131	0.080	0.290	0.288
1/1024	0.137	0.079	0.375	0.376

Test Problems - Scaling

- Standard AMG does fine on these problems
- Break assumption on local character of algebraically-smooth error
- Scale problems nodally by 10^{5r}, where r is chosen uniformly between 0 and 1 for each node

Numerical Results - Solution Time

Total time to setup MG method and then reduce residual by $10^{10}\,$

	Problem 1		Problem 2		
h	Standard AMG	Adapted AMG	Standard AMG	Adapted AMG	
1/64	*	0.03s	*	0.05s	
1/128	*	0.22s	*	0.27s	
1/256	*	0.91s	*	1.22s	
1/512	*	3.64s	*	5.35s	
1/1024	*	15.64s	*	28.27s	

Numerical Results - Solution Time

Total time to setup MG method and then reduce residual by 10^{10}

or residual reduction after 200 iterations

	Problem 1		Problem 2	
h	Standard AMG	Adapted AMG	Standard AMG	Adapted AMG
1/64	$3.3 imes 10^{-5}$	0.03s	4.5×10^{-5}	0.05s
1/128	$3.6 imes 10^{-5}$	0.22s	2.5×10^{-5}	0.27s
1/256	$2.5 imes 10^{-5}$	0.91s	1.7×10^{-5}	1.22s
1/512	1.8×10^{-5}	3.64s	1.2×10^{-5}	5.35s
1/1024	$1.3 imes 10^{-5}$	15.64s	$9.3 imes 10^{-6}$	28.27s

Numerical Results - Convergence Factors

Asymptotic convergence	factors of	of resulting	V(1,1)	cycles
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	Problem 1		Problem 2	
h	Standard AMG	Adapted AMG	Standard AMG	Adapted AMG
1/64	0.991	0.069	0.996	0.187
1/128	0.997	0.078	0.996	0.212
1/256	0.996	0.077	0.996	0.235
1/512	0.996	0.078	0.996	0.292
1/1024	0.996	0.079	0.995	0.383

Theoretical Results

Questions:

- Convergence of adaptive process
- Convergence of the resulting V-cycle
- Approached theory in a 2-level, reduction-based AMG setting
- Convergence of solution phase based on *A*-orthogonal decomposition
- Show convergence of 2-level adaptive process in reducing Rayleigh Quotient of near-null-space approximation

Summary - Upscaling

- Robust variational multigrid methods define useful coarse-scale models
- Solution of these models accurately approximates fine-scale solution
 - Net outflow flux
 - Fine-scale structure of pressure
- Accurate recovery of coarse-scale material properties
- Coarse-scale model is accurate with multiscale basis function, but includes regularization term with coarse-scale basis interpretation

Future Work - Upscaling

- Compare with multiscale finite element method (Hou et al.)
- Coarsen using AMG or adaptive AMG
- Extend to three dimensions
- Apply techniques to nonlinear problem
- Reconcile averaging theory with numerical results

Summary - Adaptive Multigrid

- Assumptions on knowledge of algebraic smoothness in classical algebraic multigrid methods can be relaxed
- Additional work in setup to expose prototypical algebraically-smooth error results in improved convergence behavior
- Prove resulting algorithm is invariant to diagonal scalings
- Theory supports convergence of adaptive process

Future Work - Adaptive Multigrid

- Apply adaptive framework to systems of PDEs
- Extend AMG interpolation to fit multiple prototypes
 - Adaptive smoothed aggregation performs well on linear elasticity
- Adaptive choice of coarse grid
- Improve and extend theory

Conclusions

- Effective multiscale basis functions created naturally in variational multigrid
- Solutions to variational coarse-scale problems accurately predict fine-scale behavior
- Adaptive process creates accurate coarse-scale models through exposure of algebraically-smooth error
- Resulting V-cycle outperforms classical AMG on many problems
- Many interesting questions remain