(Block) ILUT smoothers for *p*-multigrid methods in Isogeometric Analysis R. Tielen, M. Möller and **C. Vuik** 

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#### *p*-Multigrid methods





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# Isogeometric Analysis (IgA)

- Extension of the Finite Element Method (FEM)
- Geometry  $\Omega$  and solution *u* are approximated by same basis functions (**B-Spline basis functions**)
- Global mapping from  $\Omega_h$  to parametric domain  $\hat{\Omega}_h$
- Description of the geometry that is highly accurate  $(\Omega = \Omega_h)$  throughout all computation steps



# Model problem (CDR-equation)

Consider

$$-\nabla \cdot (\mathsf{D}\nabla u) + \mathsf{v} \cdot \nabla u + Ru = f, \quad \text{on } \Omega \quad (1)$$
$$u = g, \quad \text{on } \partial \Omega \quad (2)$$

where D denotes the diffusion tensor, v a divergence-free velocity field and R a source term. Here  $\Omega \subset \mathbb{R}^2$  is a connected, Lipschitz domain and  $f \in L^2(\Omega)$ .

Let  $\mathcal{V} = H_0^1(\Omega)$  denote the space of functions in the Sobolev space  $H^1(\Omega)$  that vanish at  $\partial\Omega$ .



#### Variational formulation

Multiplication of Equation (1) with an arbitrary test function  $v \in \mathcal{V}$  and application of integration by parts leads to the variational form:

$$a(u,v) = (f,v) \quad \forall v \in \mathcal{V},$$
 (3)

where

$$\mathsf{a}(u,v) = \int_\Omega (\mathsf{D} 
abla u) \cdot 
abla v + (\mathsf{v} \cdot 
abla u) v + \mathsf{R} u \mathsf{v} \; \mathsf{d} \Omega$$

and

$$(f,v)=\int_{\Omega} f v \ \mathrm{d}\Omega.$$



#### B-spline basis functions

Isogeometric Analysis adopts **B-spline basis functions** to discretize the variational formulation





#### B-spline basis functions

#### Properties of B-spline basis functions

- Compact support ⇒ Sparse system matrices
- Strictly positive  $\Rightarrow$  Mass matrix positive
- Partition of unity  $\Rightarrow$  Direct mass lumping



#### Galerkin formulation

Given the spline space  $\mathcal{V}_{h,p}$ , the Galerkin formulation of (3) becomes:

Find  $u_{h,p} \in \mathcal{V}_{h,p}$  such that  $a(u_{h,p}, v_{h,p}) = (f, v_{h,p}) \qquad \forall v_{h,p} \in \mathcal{V}_{h,p},$ 

where p is the approximation order of the B-splines and h the mesh width. The discretized problem can be written as a linear system

$$A_{h,p}u_{h,p} = f_{h,p}.$$
 (4)



#### Need for efficient solvers

For a fixed mesh width h, the condition number  $\kappa(A_{h,p})$  scales exponentially with the approximation order p.



#### Need for efficient solvers

#### Enhanced *h*-multigrid methods

- Subspace corrected mass smoother [Takacs, 2017]
- Hybrid smoother [Sogn, 2018]
- Multiplicative Schwarz smoother [de la Riva, 2018]

#### Preconditioners

- Schwarz methods [Beirão da Veiga, 2012]
- Sylvester equation [Sangalli, 2016]

#### Our solution strategy:

p-multigrid methods [Tielen et al, 2018 & 2020]



#### p-multigrid method

#### Motivation

The linear system A<sub>h,p</sub>u<sub>h,p</sub> = f<sub>h,p</sub>
becomes more difficult to solve for increasing p
reduces to standard C<sup>0</sup>-FEM for p = 1 (where *h*-multigrid is an established solution technique)

In contrast to *h*-multigrid methods (in IgA)

- the #DoFs remains similar on coarser *p*-levels
- the stencil reduces significantly on coarse *p*-levels

• the spaces are not nested 
$$(\mathcal{V}_{h,p} \not\supseteq \mathcal{V}_{h,p-1} \not\supseteq \dots)$$



#### p-multigrid method



#### Prolongation and restriction



Let  $\phi_i^q$  denote the *i*<sup>th</sup> basis function from  $\mathcal{V}_{h,q}$ . Then define

$$(\mathsf{M}^{r}_{q})_{(i,j)} := \int_{\hat{\Omega}_{h}} \phi^{q}_{i}(\boldsymbol{\xi}) \ \phi^{r}_{j}(\boldsymbol{\xi}) \ c(\boldsymbol{\xi}) \ \mathrm{d}\hat{\Omega}$$

Replace  $M_q^q$  by its row-sum lumped counterpart!



#### Smoother at high order level: Gauss-Seidel?



Figure: Spectrum iteration matrix *p*-multigrid for p = 2 (left) and p = 3 (right) for Poisson's equation on a quarter annulus.

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# Alternative: ILUT smoother [Saad 1994]

**Setup**: Incomplete LU factorization of  $A_{h,p} \approx L_{h,p}U_{h,p}$  thereby

• dropping all elements lower than tolerance  $\tau = 10^{-13}$ 

keeping only the N (= average number of non-zero entries in each row of A<sub>h,p</sub>) largest elements in each row

**Application**: perform  $s = 1, \ldots, \nu$  smoothing steps

$$\begin{aligned} \mathbf{e}_{h,p}^{(s)} &= (\mathsf{L}_{h,p}\mathsf{U}_{h,p})^{-1}(\mathsf{f}_{h,p}-\mathsf{A}_{h,p}\mathsf{u}_{h,p}^{(s)}) \\ \mathbf{u}_{h,p}^{(s+1)} &= \mathsf{u}_{h,p}^{(s)}+\mathsf{e}_{h,p}^{(s)} \end{aligned}$$



### Spectrum iteration matrix (p=2)



Figure: Spectrum iteration matrix for Gauss-Seidel (left) and ILUT (right) for Poisson's equation on a quarter annulus.



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### Spectrum iteration matrix (p=3)



Figure: Spectrum iteration matrix for Gauss-Seidel (left) and ILUT (right) for Poisson's equation on a quarter annulus.



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### Spectrum iteration matrix (p=4)



Figure: Spectrum iteration matrix for Gauss-Seidel (left) and ILUT (right) for Poisson's equation on a quarter annulus.



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#### Benchmark #1

• Poisson's equation:

$$\mathsf{D} = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}, \quad \mathsf{v} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}, \quad R = 0.$$

•  $\Omega$  is a quarter annulus with radia 1 and 2.

	p =	<i>p</i> = 2		<i>p</i> = 3		p = 4		p = 5	
	ILUT	GS	ILUT	GS	ILUT	GS	ILUT	GS	
$h = 2^{-6}$	4	30	3	62	3	176	3	491	
$h = 2^{-7}$	4	29	3	61	3	172	3	499	
$h = 2^{-8}$	5	30	3	60	3	163	3	473	
$h = 2^{-9}$	5	32	3	61	3	163	3	452	
	'				,		,		



#### Benchmark #2

• CDR-equation:  $D = \begin{bmatrix} 1.2 & -0.7 \\ -0.4 & 0.9 \end{bmatrix}, \quad v = \begin{bmatrix} 0.4 \\ -0.2 \end{bmatrix}, \quad R = 0.3.$ 

•  $\Omega$  is the unit square, i.e.  $\Omega = [0, 1]^2$ .

	<i>p</i> = 2		<i>p</i> = 3		<i>p</i> = 4		<i>p</i> = 5	
	ILUT	GS	ILUT	GS	ILUT	GS	ILUT	GS
$h = 2^{-6}$	5	-	3	_	3	_	4	_
$h = 2^{-7}$	5	-	3	_	4	_	4	_
$h = 2^{-8}$	5	_	3	_	3	_	4	_
$h = 2^{-9}$	5	_	4	_	3	_	4	_



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# CPU timings (single solve)

- Comparison with *h*-multigrid method [Takacs,2017]
- Higher setup costs with *p*-multigrid, but fast solves!



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# CPU timings (Multiple solves)



• Increasing influence of solving costs



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#### Multipatch geometries

- Geometry Ω<sub>h</sub> can **not always** be described by a single mapping to parametric domain Ω<sub>h</sub>!
- Represent Ω by non-overlapping subdomains (patches), each with their own mapping
- Resulting operator A<sub>h,p</sub> has a block structure, where each patch leads to a single block



#### Multipatch geometries



Figure: Sparsity pattern of the system matrix (left) and global ILUT factorization (right) for p = 3 and  $h = 2^{-5}$  for Poisson's equation on a quarter annulus.



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# Block ILUT

#### Observation

We can write A = LU as:

$$\begin{bmatrix} A_{11} & 0 & A_{\Gamma 1} \\ & \ddots & & \vdots \\ 0 & A_{KK} & A_{\Gamma K} \\ A_{1\Gamma} & \cdots & A_{K\Gamma} & A_{\Gamma\Gamma} \end{bmatrix} = \begin{bmatrix} L_1 & & \\ & \ddots & \\ B_1 & \cdots & B_K & I \end{bmatrix} \begin{bmatrix} U_1 & & C_1 \\ & \ddots & & \vdots \\ & & U_K & C_K \\ & & & S \end{bmatrix}$$

where

• 
$$A_{ii} = L_i U_i$$
  
•  $B_i = A_{i\Gamma} U_i^{-1}$   
•  $C_i = L_i^{-1} A_{\Gamma i}$   
•  $S = A_{\Gamma\Gamma} - \sum_{i=1}^{K} B_i C_i$ 



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# Block ILUT

#### Key Idea (Nievinski 2018)

Replace L and U by their ILUT factorizations:

$$\begin{bmatrix} A_{11} & A_{\Gamma 1} \\ & \ddots & \vdots \\ & A_{KK} & A_{\Gamma K} \\ A_{1\Gamma} & \cdots & A_{K\Gamma} & A_{\Gamma\Gamma} \end{bmatrix} \approx \begin{bmatrix} \tilde{L}_{1} & & & \\ & \ddots & & \\ \tilde{B}_{1} & \cdots & \tilde{B}_{K} & I \end{bmatrix} \begin{bmatrix} \tilde{U}_{1} & \tilde{C}_{1} \\ & \ddots & \vdots \\ & \tilde{U}_{K} & \tilde{C}_{K} \\ & \tilde{S} \end{bmatrix},$$
  
where  
$$\textcircled{A}_{ii} \stackrel{(!)}{\approx} \tilde{L}_{i} \tilde{U}_{i}$$
$$\textcircled{B}_{i} = A_{i\Gamma} \tilde{U}_{i}^{-1}$$
$$\textcircled{B}_{i} = A_{i\Gamma} \tilde{U}_{i}^{-1}$$
$$\textcircled{S} = A_{\Gamma\Gamma} - \sum_{i=1}^{K} \tilde{B}_{i} \tilde{C}_{i}$$

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### Block ILUT

•  $\tilde{L}_i, \tilde{U}_i$  can be determined in parallel

• Inversion of  $\tilde{L}_i, \tilde{U}_i$  is avoided by solving:

$$\tilde{\mathsf{U}}_i^{\top}\tilde{\mathsf{B}}_i^{\top}=\mathsf{A}_{i\Gamma}^{\top},\qquad \tilde{\mathsf{L}}_i\tilde{\mathsf{C}}_i=\mathsf{A}_{\Gamma i},$$



Figure: Global and block ILUT factorization for Poisson's equation

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### Block ILUT vs. ILUT

#### Poisson on Yeti Footprint

	p = 2		<i>p</i> = 3		p =	= 4	p = 5	
	Global	Block	Global	Block	Global	Block	Global	Block
$h = 2^{-3}$	5	4	4	2	4	2	4	2
$h = 2^{-4}$	8	4	5	3	5	3	4	2
$h = 2^{-5}$	8	4	6	3	5	3	5	3





#### Conclusions



are efficient and robust solvers for IgA

enhanced with ILUT as a smoother they are

- robust in the order p and mesh width h
- competitive to state-of-the-art h-multigrid methods

 adopting block ILUT has potential (for parallelization) in case of multipatch geometries.

#### Further reading:

R.Tielen, M. Möller, D. Göddeke and C.Vuik: p-multigrid methods and their comparison to h-multigrid methods within Isogeometric Analysis, Comput. Methods Appl. Mech. Engrg., Vol 372 (2020)





### Outlook

- Theoretical insight into effectiveness of ILUT
- Further exploration of block ILUT smoother
- Exploit parallelism of block ILUT smoother



G+Smo (Geometry plus Simulation modules), http://github.com/gismo



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Thank you for your attention! Questions?



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