

Solution of the vector wave equation using a Krylov solver with an algebraic multigrid approximated preconditioner

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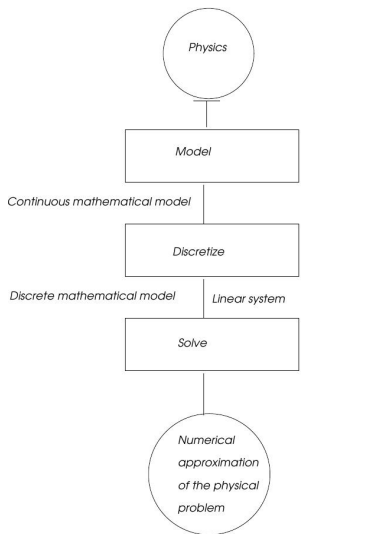
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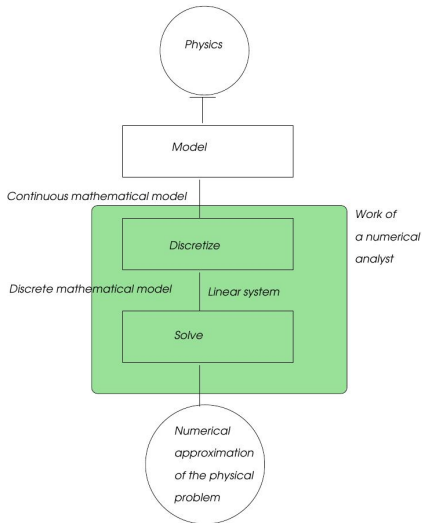
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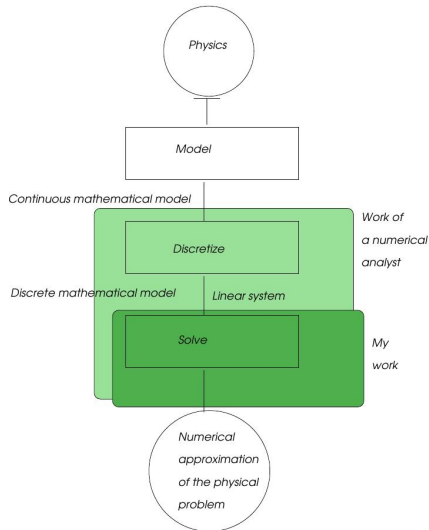
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- **RADAR (Radio Detection and Ranging):** technology to detect military platforms (e.g. aircraft, ship or tank) by using electromagnetic waves.
- **Goal:** identify the range, altitude, direction, or speed of both moving and fixed objects.
- **Measure of detectability:** **radar cross section (RCS)** (depends on observation angle, frequency, polarization).
- A platform is detected when the *received* signal-to-noise ratio exceeds a certain threshold.

Problem description

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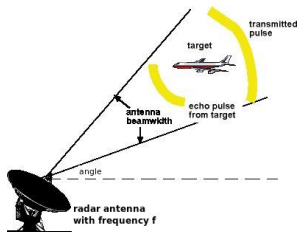
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- When the radar signature of a military platform cannot be determined experimentally, *numerical prediction* techniques are used.
- In this thesis a *faster* and more *memory efficient* prediction method is proposed.

Importance of inlet cavity scattering

Radar excited from the front: jet engine air *intake* (of a modern fighter aircraft) accounts for the major part of the RCS for a large range of observation angles.

Figure: Jet engine air intake (left) closed by jet engine compressor fan (right) – together forms a **large and deep open-ended cavity** with varying cross section – *Dimensions: $d \approx 30\lambda$ and $L \approx 200\lambda$ for X-band excitation (10 GHz).*



General Maxwell equations – Assumptions

- Start with general Maxwell equations: describe properties of electric (\mathbf{E}) and magnetic (\mathbf{H}) fields and relate them to their source: an electric current density distribution.
- Assume that the field quantities are harmonic oscillating functions with an angular frequency ω^* (*time-harmonic functions*).

And He said

$$\begin{aligned} \oint \bar{\mathbf{E}} \cdot d\bar{\mathbf{s}} &= - \int \frac{\partial \bar{\mathbf{B}}}{\partial \tau} \cdot d\bar{\mathbf{s}} & \nabla \times \bar{\mathbf{E}} &= -\mu \frac{\partial \bar{\mathbf{H}}}{\partial \tau} & \nabla \times \bar{\mathbf{E}} &= -\mu \frac{\partial \bar{\mathbf{H}}}{\partial \tau} \\ \oint \bar{\mathbf{H}} \cdot d\bar{\mathbf{s}} &= \int \left(\bar{\mathbf{J}}_c + \frac{\partial \bar{\mathbf{D}}}{\partial \tau} \right) \cdot d\bar{\mathbf{s}} & \text{OR } \nabla \times \bar{\mathbf{H}} &= \bar{\mathbf{J}}_c + \varepsilon \frac{\partial \bar{\mathbf{E}}}{\partial \tau} & \text{OR } \nabla \times \bar{\mathbf{H}} &= \bar{\mathbf{J}}_c + \varepsilon \frac{\partial \bar{\mathbf{E}}}{\partial \tau} \\ \oint \bar{\mathbf{D}} \cdot d\bar{\mathbf{s}} &= \int \bar{\mathbf{v}} \cdot \bar{\mathbf{D}} d\bar{v} & \nabla \cdot \bar{\mathbf{D}} &= \rho_c & \nabla \cdot \bar{\mathbf{D}} &= \rho_c \\ \oint \bar{\mathbf{B}} \cdot d\bar{\mathbf{s}} &= 0 & \nabla \cdot \bar{\mathbf{B}} &= 0 & \nabla \cdot \bar{\mathbf{B}} &= 0 \end{aligned}$$

and there was light

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Use the *constitutive relations*¹, to derive the dimensionless *vector wave equation* for the electric field:

$$\nabla \times \nabla \times \mathbf{E} - k_0^2 \varepsilon_r \mathbf{E} = -jk_0 Z_0 \mathbf{J}. \quad (1)$$

(same type of equation for the magnetic field \mathbf{H} .)

Boundary conditions to obtain a well-posed problem:

- Impose *global* absorbing boundary conditions on the *aperture*: they lead to more accurate numerical solutions.
- Impose Dirichlet boundary conditions on the cavity surface.

¹Used to describe the material properties of the medium of interest.

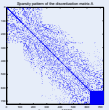
FEM discretization

Equation (1) is discretized by a higher order edge based finite element discretization method, resulting in a **large** linear system:

$$Ax = b.$$

Properties system matrix A :

- Consists of a sparse part and a fully populated part,
- ‘Nearly’ symmetric but not Hermitian,
- Ill-conditioned,
- Has unfavourable spectrum.



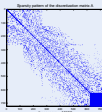
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- Use a preconditioner matrix M .
- Solve $M^{-1} A x = M^{-1} b$.

Based on the work of Erlangga²:

- **Bi-CGSTAB** method to solve the preconditioned system:
`bicgstab(A, b, ITER-MAX, TOL, M)`.
- Apply *shifted Laplace preconditioner* and solve preconditioner system with multigrid.
- Helmholtz and Maxwell's equation have similar properties.
⇒ Expected that the shifted Laplace preconditioner and multigrid will also be very effective to incorporate in the current application.

²See my report for the complete reference.

Homogeneous Helmholtz equation

- $-\Delta u - k_0^2 u = 0.$
- Shifted Laplace operator: $\mathcal{M}_{(\beta_1, \beta_2)} = -\Delta - (\beta_1 + \iota\beta_2)k_0^2.$

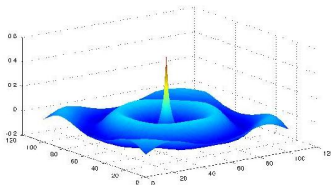
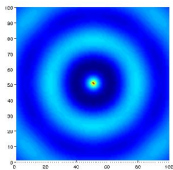
Vector wave equation

- Vector form of the Helmholtz equation:
 $-\Delta \mathbf{E} - k_0^2 \mathbf{E} = 0, \text{ with } \mathbf{E} = (E_x, E_y, E_z).$
- Shifted Laplace operator in *vector* form:
 $\mathcal{W}_{(\beta_1, \beta_2)} = -\Delta - (\beta_1 + \iota\beta_2)k_0^2.$

Solve the Helmholtz equation using Bi-CGSTAB–AMG

Multigrid type

- Erlangga used *geometric* multigrid to perform the preconditioner solve: Bi-CGSTAB–*MG* algorithm.
- In this experiment *algebraic* multigrid is used with one V-cycle, leading to the same solution as Erlangga.



Surface plot of the real part of the solution of the *two dimensional* Helmholtz equation with local absorbing boundary conditions.

Present implementation

- Dimension of system matrix $A : N \approx 1 \cdot 10^7$ (total number of unknowns).
- Iterative Krylov subspace method: *Generalized Conjugate Residual* (GCR) method.
- Ill-conditionedness and unfavourable spectrum of A
 \Rightarrow improve convergence GCR by *shifted Laplace preconditioner*.
- GCR long recurrence method \Rightarrow difficult to satisfy memory requirements for storing eigenvectors of Krylov basis.

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Proposed solution:

Modify existing algorithm by using a *multigrid solution method* for the shifted Laplace preconditioner system and use a *short recurrence method* for the preconditioned system.

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Description of the methods

- IDR(4) short recurrence Krylov method from Van Gijzen and Sonneveld.
- ML: Sandia's laboratories main multigrid preconditioning package³.

³See my report for the complete references.

Description of the methods

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Limitations using ML

- ML can only perform computations in real valued arithmetic.
- Use optimal **red** shift for the shifted Laplace preconditioner.
- Only applications with non-absorbing materials (permittivity and permeability) can be considered.

³See my report for the complete references.

IDR(4)–ML-AMG algorithm for a small cavity scattering model with dimensions $1.5\lambda \times 1.5\lambda \times 0.6\lambda$

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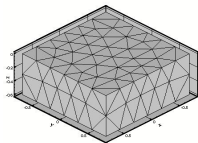
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p	h	N	$(-1, 0)$	$(1, -0.5)$	ML-solve $(-1, 0)$
0	0.25	1402	245	150	308
0	0.20	2796	301	226	374
1	0.35	2914	343	270	394
1	0.30	4344	342	427	417
1	0.25	7960	413	459	530
2	0.35	5316	390	477	546
2	0.30	8730	473	494	833

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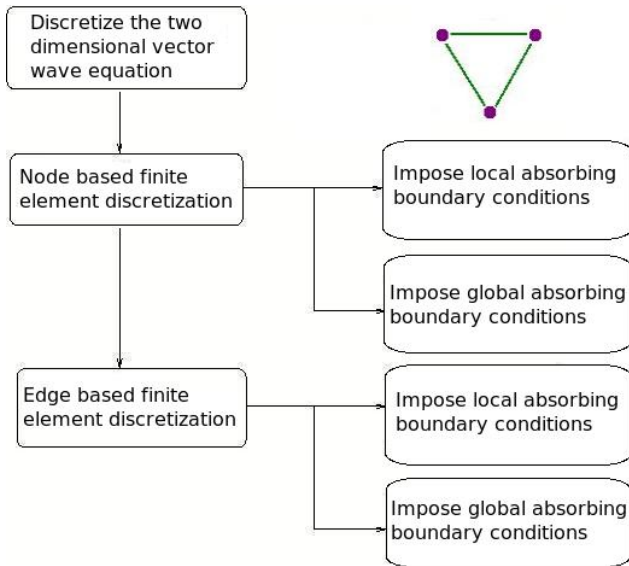
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To improve the performance of the IDR(4)–ML-AMG algorithm, the two dimensional vector wave equation is considered \Rightarrow two dimensional Maxwell solver⁴.

- Relatively small number of unknowns for high wavenumbers.
- Extend results to three dimensional case.

⁴Made by Duncan van der Heul and Shiraz Abdoel.

Flowchart two dimensional vector wave equation



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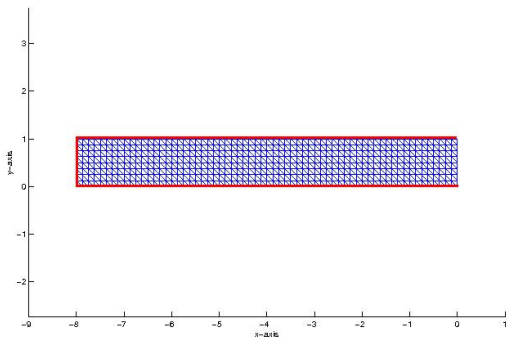


Figure: Mesh of a two dimensional cavity with dimensions 8×1 .

Node based implementation: E_z solution

The equation considered here:

$$-\Delta E_z - k_0^2 \varepsilon_r E_z = 0$$

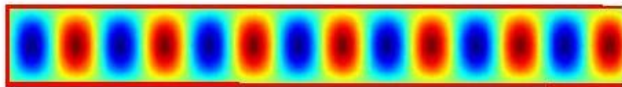


Figure: Contour plot of the real part of E_z . It can be seen how the wave travels through the inlet and is reflected on the bottom.

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- Test cases with varying wavenumber and varying mesh size: IDR(4)-ML-AMG algorithm performs good.
- Several test cases with different boundary conditions.

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- Several test cases with different boundary conditions.

k_0	2π	4π	6π
(A_{loc}, M_{loc})	83	225	440
(A_{gl}, M_{loc})	79	247	379
(A_{gl}, M_{gl})	75	239	427

Table: Total number of matrix vector operations for the IDR(4)-ML-AMG algorithm.

The equation considered here:

$$\nabla \times \nabla \times \mathbf{E} - k_0^2 \varepsilon_r \mathbf{E} = 0, \text{ with } \mathbf{E} = (E_x, E_y, 0)^T.$$

$$\text{It holds that: } \nabla \times \mathbf{E} = i\omega \mathbf{H} = i\omega(0, 0, H_z)^T.$$

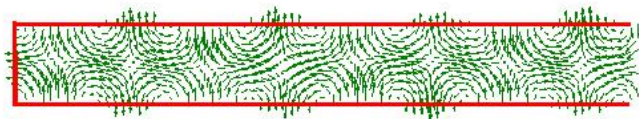


Figure: Vector plot of E_x and E_y .

Edge based implementation: H_z solution

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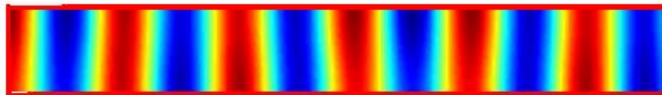


Figure: Contour plot of the real part of H_z .

Edge based implementation with local ABC similar to node based implementation.

Conclusions for two dimensional case

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- Work involved in exact solve preconditioner system versus single V-cycle AMG solve similar for *real* shift \Rightarrow Expect the same for optimal *complex* shift.
- Use preconditioner with operator based on *local* ABC \Rightarrow more efficiently solved (e.g. with a complex algebraic multigrid method).
- Optimal real shift is a *restriction* compared to optimal complex shift.

Three dimensional cavity

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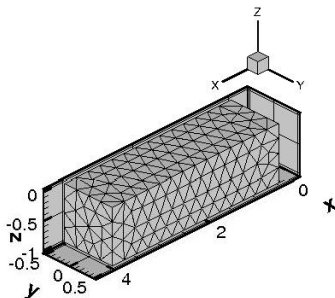
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mesh size h	0.10
degrees of freedom N	79,428
wave number k_0	2π
wave length λ	1
preconditioner system	solved using ML for optimal real shift exact solve for optimal complex shift

Performance of nested GCR algorithm versus the IDR(4)–ML-AMG algorithm

Nested GCR algorithm

- Optimal shift used: $(\beta_1, \beta_2) = (0.5, 3.0)$.
- Orthogonalisation of all basisvectors \Rightarrow lot of work.
- # MAT-VEC-OPs: 423,000.

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- # MAT-VEC-OPs: 423,000.

Algorithm	$(\beta_1, \beta_2) = (1, -0.5)$	$(\beta_1, \beta_2) = (-1, 0)$
IDR(4)	543	–
IDR(4)–ML-AMG	–	5307
expected factor	800	80

Table: Total number of matrix vector operations for the different algorithms and the expected factor which can be gained compared to the currently used algorithm.

- Restriction to real valued arithmetic leads to a reduction in work by factor 100.
- Optimal real shift restriction on effectiveness which can be gained using optimal complex shift.
- Optimal complex shift leads to a reduction in work by factor 1000.

Recommendations for future research

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- Incorporate IDR(s) method in the existing algorithm.
- Introduce preconditioner based on local absorbing boundary conditions.

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Questions & Discussion