Midterm review:

Mimetic

Isogemetric FEM

M.Sc. Thesis project

by Stevie-Ray Janssen









Combine ideas from isogeometric analysis and mimetic methods to develop a structure-preserving discretization for the Euler equations for incompressible fluids.



Project outline

• Planning:

Start Date:	1-9-2015	dinsdag	-																																									
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Nearing/Orienting	9-1-15	9-10-15	10	100%	0	10	0																																					
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Fundamentals Phase I	10-1-15	1-3-16	95	100%	67	95	0																																					
Reading	10-1-15	10-27-15	27	90%	19	24	3																																					
Work, Implementation, and Analysis	10-28-15	11-23-15	27	90%	19	24	3						-																															
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Holiday	12-21-15	1-3-16	14	100%	10	14	0													i –																								
Euler equations, Phase II	1-4-16	3-31-16	88	30%	64	26	62												_																									
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Iterative solving, Phase III	4-1-16	6-26-16	87	0%	61	0	87																																					
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Hand in Literature Report	6-26-16	6-26-16	1	0%	0	0	1																																					
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Verification/Validation of Results	7-9-16	7-21-16	13	0%	9	0	13																																	۰.				
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Project outline (cnt'd)

- Phase I questions:
 - How can we use IGA to solve PDE's?
 - What structures are facilitated in elliptic PDE's?
 - How can we preserve these structures?
 - Can we construct a MIMIGA method to discretize an elliptic PDE problem?



This presentation - literature review

- Introduction
 - Isogeometric Analysis & Mimetic Methods
- Approach for elliptic PDE's
 - Exterior calculus
 - DeRham complex
 - Application: Scalar Poisson equation in 2D
- Conclusion
- Future work



Introduction – Isogeometric Analysis

- Introduced by the Hughes group in 2005 to bridge the gap between CAD and FEM
- Isogeometric paradigm



B-splines make an excellent basis for FEM





Introduction – Mimetic Methods

- PDE's facilitate physical structures and symmetries.
- Tools from exterior calculus and algebraic topology are used to capture these structures.
- Growing awareness: Disrete exterior calculus, discrete hodge theory, exterior finite element method, compatible methods, mimetic finite diference, etc



Why exterior calculus?

- Structures become apparent.
- Distinction between topological and metric dependencies.
- Generalized for *n* dimensions.



Differential Forms; $\alpha^{(k)}$

- Differential forms are elements from the dual vector space,
- Associated with geometric structure,

- 0-form:
$$f^{(0)} = f(x, y)$$

- 1-form: $\alpha^{(1)} = \alpha_1(x, y)dx + \alpha_2(x, y)dy$
- "Measurement of physical variables,"

$$-M = \oiint \rho^{(2)} = \oiint \rho(x, y) dx \wedge dy$$

Space of k-forms: Λ^(k)



Exterior derivative; d

- Exterior derivative d generalizes ∇f , $\nabla \times \underline{\omega}$, $\nabla \cdot \underline{v}$ $d\alpha^{(1)} = \left(\frac{\partial \alpha_2}{\partial x} - \frac{\partial \alpha_1}{\partial y}\right) dx \wedge dy$ • $d: \Lambda^{(k)} \to \Lambda^{(k+1)}$ $\mathbb{R} \longrightarrow \underbrace{\left(\frac{1}{2} - \frac{1}{2}\right)^2}_{\Lambda^{(0)} = \frac{1}{2} - \frac{1}{2}} \xrightarrow{d} \underbrace{\left(\frac{1}{2} - \frac{1}{2}\right)^2}_{\Lambda^{(1)} = \frac{1}{2} - \frac{1}{2}} \xrightarrow{d} \underbrace{\left(\frac{1}{2} - \frac{1}{2}\right)^2}_{\Lambda^{(2)} = \frac{1}{2} - \frac{1}{2}} \xrightarrow{d} \underbrace{\left(\frac{1}{2} - \frac{1}{2}\right)^2}_{\Lambda^{(2)} = \frac{1}{2} - \frac{1}{2}} \xrightarrow{d} \underbrace{\left(\frac{1}{2} - \frac{1}{2}\right)^2}_{\Lambda^{(2)} = \frac{1}{2} - \frac{$
- Exact sequence, the DeRham complex
- Nilpotent, $dd\alpha^{(k)} = 0$
- Independent of metric

Hodge-* operator;

- Maps forms to dual geometry,
- Metric dependent,
- Double DeRham complex,





Codifferential; d^*

- $d^* \coloneqq d \star$
- Adjoint of d: $(\cdot, d^* \cdot) = (d \cdot, \cdot) \int bc's$
- Laplace operator: $\Delta = dd^* + d^*d$





Scalar Poisson equation

- E.g. Potential flow, electrostatics,
- Given $f(x, y) = 2\pi^2 \sin(\pi x) \sin(\pi y)^{-1} \int_{\lambda_{2}^{1/2}}^{0.4} \int_{0.8}^{0.4} \int_{10^{0.0}}^{0.0} find \varphi(x, y)$ such that $\Delta \varphi = f$ on $\Omega = [0,1]^2$ with $\varphi = 0$ on $\partial \Omega$

0-form,

2-form,

Find $\varphi^{(0)}$ s.t. $d^*d\varphi^{(0)} = f^{(0)}$

Find $\sigma^{(2)}$ s.t. $dd^* \sigma^{(2)} = f^{(2)}$

Same solution, different discretization



1.0 0.8

- 0.6 - 0.4 u[-] - 0.2 - 0.0 - 0.2 - 0.4

< 1.0 0.8

0.6 0.4 2

0.2

0-form Poisson; $d^*d\varphi^{(0)} = f^{(0)}$

- Weak formulation, $\begin{pmatrix} w^{(0)}, d^*d\varphi^{(0)} \end{pmatrix}_{\Omega} = \begin{pmatrix} w^{(0)}, f^{(0)} \end{pmatrix}_{\Omega} \\ \Leftrightarrow \\ \begin{pmatrix} dw^{(0)}, d\varphi^{(0)} \end{pmatrix}_{\Omega} = \begin{pmatrix} w^{(0)}, f^{(0)} \end{pmatrix}_{\Omega} - \oint_{\partial\Omega} w^{(0)} \wedge \star d\varphi^{(0)} \end{pmatrix}_{\Omega}$
- Well-posedness through Lax-Milgram,



0-form Poisson; FEM

- Conforming FEM, take $\Lambda_{h}^{(k)} \subset \Lambda^{(k)}$
- Use B-spline spans $\Lambda_{h}^{(0)} = S^{p,p}$





0-form Poisson; edge functions

• Applying the exterior derivative (1D-example)

– Nodal basis:
$$\varphi_h^{(0)} = \sum_{i=0}^n \varphi_i h_i^p(x) = \left(\underline{\varphi}\right)^T \underline{R}^0$$

- Then,
$$d\varphi_h^{(0)} = \sum_{i=1}^n (\varphi_i - \varphi_{i-1}) e_i^{p-1}(x) = \left(\mathbb{E}^{(10)}\underline{\varphi}\right)^T \underline{R}^1$$

Differences of coefficients are captured in matrix using {-1,0,1}

New edge type basis function emerges with a polynomial degree less



0-form Poisson; edge functions (cnt'd)

- Extension to 2D using tensor products of nodal and edge type basis
- Nodal/edge
 - 0-form
 - 1-form
 - 2-form



0-form Poisson, Matrices

•
$$\left(dw_{h}^{(0)}, d\varphi_{h}^{(0)}\right)_{\Omega} = \underline{w}^{T} (\mathbb{E}^{10})^{T} \left(\int_{\Omega} \left(\underline{R}^{(1)}\right)^{T} \underline{R}^{(1)}\right) (\mathbb{E}^{10}) \underline{\varphi}$$

• Result: $\left(\mathbb{E}^{(10)}\right)^T \mathbb{M}^{(11)} \mathbb{E}^{(10)} \varphi = f$

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0-form Poisson, Matrices (cnt'd)

- Exact discretization of $v^{(1)} = d\varphi^{(0)}$ through incidence matrices, $\underline{v} = \mathbb{E}^{(10)}\varphi$
- Incidence matrices are nilpotent $\mathbb{E}^{(21)}\mathbb{E}^{(10)} = \emptyset$, and satisfy the DeRham sequence
- Hodge-* operator (metric) is discretized through mass matrix M⁽¹¹⁾





0-form Poisson, Results



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0-form Poisson, Results (cnt'd)



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2-form Poisson; $dd^*\sigma^{(2)} = f^{(2)}$

Weak formulation;

$$\left(w^{(2)},dd^{*}\sigma^{(2)}\right)_{\Omega}=\left(w^{(2)},f^{(2)}\right)_{\Omega}$$

Integration by parts? No, take mixed formulation:

$$\begin{cases} d^* \sigma^{(2)} = \psi^{(1)} \\ d\psi^{(1)} = f^{(2)} \end{cases}$$

• Weak form:

$$\begin{cases} \left(dq^{(1)}, \sigma^{(2)}\right)_{\Omega} = \left(q^{(1)}, \psi^{(1)}\right)_{\Omega} - \oint_{\partial\Omega} q^{(1)} \wedge \star \sigma^{(2)} \\ \left(w^{(2)}, d\psi^{(1)}\right)_{\Omega} = \left(w^{(2)}, f^{(2)}\right)_{\Omega} \end{cases}$$

Well posedness through Inf-Sup conditions

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2-form Poisson; FEM

Can we take,

$$-\Lambda_{h}^{(1)} = S^{p,p}?$$
$$-\Lambda_{h}^{(2)} = S^{p,p}?$$

 No, well-posedness depends on the DeRham sequence. We take

$$-\Lambda_{h}^{(1)} = S^{p-1,p} \times S^{p,p-1}$$

$$-\Lambda_h^{(2)} = S^{p-1,p-1}$$

• Which satisfy exact sequence

$$S^{p,p} \xrightarrow[\mathsf{d}]{\mathbb{E}^{(10)}} S^{p-1,p} \times S^{p,p-1} \xrightarrow[\mathsf{d}]{\mathbb{E}^{(21)}} S^{p-1,p-1}$$

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• Or
$$\left(\mathbb{M}^{(22)}\mathbb{E}^{(21)}\right)^T \left(\mathbb{M}^{(11)}\right)^{-1} \left(\mathbb{M}^{(22)}\mathbb{E}^{(21)}\right) \underline{\Psi} = \underline{f}$$

$$\begin{bmatrix} -\mathbb{M}^{(11)} & \left(\mathbb{M}^{(22)}\mathbb{E}^{(21)}\right)^T \\ \mathbb{M}^{(22)}\mathbb{E}^{(21)} & \emptyset \end{bmatrix} \begin{bmatrix} \underline{\psi} \\ \underline{\sigma} \end{bmatrix} = \begin{bmatrix} \underline{0} \\ \underline{f} \end{bmatrix}$$

$$\begin{cases} -(q^{(1)}, \psi^{(1)})_{\Omega} + (dq^{(1)}, \sigma^{(2)})_{\Omega} = 0 \\ (w^{(2)}, d\psi^{(1)})_{\Omega} = (w^{(2)}, f^{(2)})_{\Omega} \end{cases}$$

$$\begin{cases} -(q^{(1)}, \psi^{(1)})_{\Omega} + (dq^{(1)}, \sigma^{(2)})_{\Omega} = 0\\ (w^{(2)}, d\psi^{(1)})_{\Omega} = (w^{(2)}, f^{(2)})_{\Omega} \end{cases}$$

2-form Poisson; Results



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Conclusion

- Elliptic problems can be discretized using mass matrices and incidence matrices.
- Solution spaces are chosen such that they satisfy the DeRham complex.



Conclusion (cnt'd)

• Comparison 0-form & 2-form Poisson:

0-form	2-form									
$\left(\mathbb{E}^{(10)} ight)^T\mathbb{M}^{(11)}\mathbb{E}^{(10)}$	$\begin{bmatrix} -\mathbb{M}^{(11)} & \left(\mathbb{M}^{(22)}\mathbb{E}^{(21)}\right)^T \\ \mathbb{M}^{(22)}\mathbb{E}^{(21)} & \emptyset \end{bmatrix}$									
Obtain solution $arphi^{(0)}$	Obtain solutions $\sigma^{(2)}$, $\psi^{(1)}$									
Dirichlet is essentialNeumann is natural	Dirichlet is naturalNeumann is essential									
Gradient exact $\mathbb{E}^{(10)}\underline{\varphi} = \underline{v}$	Divergence exact $\mathbb{E}^{(21)} \underline{\Psi} = \underline{0}$ i.e. $\nabla \cdot v = 0$									



Future Work

- Towards the incompressible Euler equations:
 - Extend to hyperbolic problems,
 - Linear advection equation.
 - Construction of periodic domain.



- Staggering velocity and vorticity in time?

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Questions

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