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ANSWERS OF THE TEST NUMERICAL METHODS FOR DIFFERENTIAL EQUATIONS (CTB2400 WI3097TU) Thursday August 11th 2016, 18:30-21:30

1. (a) The local truncation error is defined as

$$\tau_{n+1}(\Delta t) = \frac{y_{n+1} - z_{n+1}}{\Delta t},$$
(1)

where z_{n+1} is given by

$$z_{n+1} = y_n + \Delta t \left(a_1 f(t_n, y_n) + a_2 f(t_n + \Delta t, y_n + \Delta t f(t_n, y_n)) \right).$$
(2)

A Taylor expansion of f around (t_n, y_n) yields

$$f(t_n + \Delta t, y_n + \Delta t f(t_n, y_n)) = f(t_n, y_n) + \Delta t \frac{\partial f}{\partial t}(t_n, y_n) + \Delta t f(t_n, y_n) \frac{\partial f}{\partial y}(t_n, y_n) + O((\Delta t)^2).$$
(3)

This is substituted into equation (2) to obtain

$$z_{n+1} = y_n + \Delta t \left(a_1 f(t_n, y_n) + a_2 \left[f(t_n, y_n) + \Delta t \frac{\partial f}{\partial t}(t_n, y_n) + \Delta t f(t_n, y_n) \frac{\partial f}{\partial y}(t_n, y_n) \right] \right) + O((\Delta t)^3$$
(4)

A Taylor series for y(t) around t_n gives for y_{n+1}

$$y_{n+1} = y(t_n + \Delta t) = y_n + \Delta t y'(t_n) + \frac{(\Delta t)^2}{2} y''(t_n) + O((\Delta t)^3).$$
(5)

From the differential equation we know that:

$$y'(t_n) = f(t_n, y_n) \tag{6}$$

From the Chain Rule of Differentiation, we derive

$$y''(t_n) = \frac{df(t_n, y_n)}{dt} = \frac{\partial f(t_n, y_n)}{\partial t} + \frac{\partial f(t_n, y_n)}{\partial y}y'(t_n)$$
(7)

after substitution of the differential equation one obtains:

$$y''(t_n) = \frac{\partial f(t_n, y_n)}{\partial t} + \frac{\partial f(t_n, y_n)}{\partial y} f(t_n, y_n)$$
(8)

Equations (5) and (4) are substituted into relation (1) to obtain

$$\tau_{n+1}(\Delta t) = f(t_n, y_n)(1 - (a_1 + a_2)) + \Delta t \left(\frac{\partial f}{\partial t} + f \frac{\partial f}{\partial y}\right) \left(\frac{1}{2} - a_2\right) + O((\Delta t)^2)$$
(9)

Hence

i. $a_1 + a_2 = 1$ implies $\tau_{n+1}(\Delta t) = O(\Delta t);$

ii. $a_1 + a_2 = 1$ and $a_2 = 1/2$, that is, $a_1 = a_2 = 1/2$, gives $\tau_{n+1}(\Delta t) = O((\Delta t)^2)$.

(b) The test equation is given by

$$y' = \lambda y. \tag{10}$$

Application of the predictor step to the test equation gives

$$w_{n+1}^* = w_n + \lambda \Delta t w_n = (1 + \lambda \Delta t) w_n.$$
(11)

The corrector step yields

$$w_{n+1} = w_n + \Delta t \left(a_1 \lambda w_n + a_2 \lambda (1 + \lambda \Delta t) w_n \right) = (1 + (a_1 + a_2) \lambda \Delta t + a_2 (\lambda \Delta t)^2) w_n.$$
(12)

Hence the amplification factor is given by

$$Q(\lambda \Delta t) = 1 + (a_1 + a_2)\lambda \Delta t + a_2(\lambda \Delta t)^2.$$
(13)

(c) Let $\lambda < 0$ (so λ is real), then, for stability, the amplification factor must satisfy

$$-1 \le Q(\lambda \Delta t) \le 1,\tag{14}$$

from the previous assignment, we have

$$-1 \le 1 + (a_1 + a_2)\lambda\Delta t + a_2(\lambda\Delta t)^2 \le 1 \Leftrightarrow -2 \le (a_1 + a_2)\lambda\Delta t + a_2(\lambda\Delta t)^2 \le 0.$$
(15)

First, we consider the left inequality:

$$a_2(\lambda\Delta t)^2 + (a_1 + a_2)\lambda\Delta t + 2 \ge 0 \tag{16}$$

For $\lambda \Delta t = 0$, the above inequality is satisfied. The discriminant of the quadratic equation is given by $(a_1 + a_2)^2 - 8a_2$. From the given assumption it follows that $(a_1 + a_2)^2 - 8a_2 < 0$ so the quadratic equation does not have real roots. Hence the left inequality in relation (15) is always satisfied. Next we consider the right hand inequality of relation (15)

$$a_2(\lambda\Delta t)^2 + (a_1 + a_2)\lambda\Delta t \le 0.$$
(17)

This relation is rearranged into

$$a_2(\lambda \Delta t)^2 \le -(a_1 + a_2)\lambda \Delta t, \tag{18}$$

hence

$$a_2|\lambda\Delta t|^2 \le (a_1+a_2)|\lambda\Delta t| \Leftrightarrow |\lambda\Delta t| \le \frac{a_1+a_2}{a_2}, \qquad a_2 \ne 0.$$
 (19)

This results into the following condition for stability

$$\Delta t \le \frac{a_1 + a_2}{a_2 |\lambda|}, \qquad a_2 \ne 0. \tag{20}$$

(d) In order to compute the Jacobian, we note that the right-hand side of the non linear system is given by:

$$f_1(x_1, x_2) = -\sin x_1 + 2x_2 + t$$
$$f_2(x_1, x_2) = x_1 - x_2^2$$

From the definition of the Jacobian it follows that:

$$\begin{pmatrix} \frac{\partial f_1}{\partial x_1} & \frac{\partial f_1}{\partial x_2}\\ \frac{\partial f_2}{\partial x_1} & \frac{\partial f_2}{\partial x_2} \end{pmatrix} = \begin{pmatrix} -\cos x_1 & 2\\ 1 & -2x_2 \end{pmatrix}$$

Substitution of $\begin{pmatrix} x_1(0) \\ x_2(0) \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$ shows that

$$J = \begin{pmatrix} -1 & 2\\ 1 & -2 \end{pmatrix}.$$

(e) For the stability it is sufficient to check that $|Q(\lambda_i \Delta t)| \leq 1$ for all the eigenvalues of the Jacobian matrix. It is easy to see that the eigenvalues of the Jacobian matrix are $\lambda_1 = -3$ and $\lambda_2 = 0$.

For the choice $a_1 = a_2 = \frac{1}{2}$ we use the expression

$$Q(\lambda \Delta t) = 1 + \lambda \Delta t + \frac{1}{2} (\lambda \Delta t)^2$$

For $\lambda_2 = 0$ it appears that $Q(\lambda_2 \Delta t) = 1$ so the inequality is satisfied for all Δt . For $\lambda_1 = -3$ we have to check the following inequalities:

$$-1 \le 1 - 3\Delta t + \frac{9}{2}(\Delta t)^2 \le 1$$

For the left-hand inequality we arrive at

$$0 \le \frac{9}{2}(\Delta t)^2 - 3\Delta t + 2$$

It appears that the discriminant $9 - 4 \cdot \frac{9}{2} \cdot 2$ is negative, so there are no real roots which implies that the inequality is satisfied for all Δt .

For the right-hand inequality we get

$$-3\Delta t + \frac{9}{2}(\Delta t)^2 \le 0$$
$$\frac{9}{2}(\Delta t)^2 \le 3\Delta t$$

 \mathbf{SO}

$$\Delta t \le \frac{2}{3}$$

(another option is to see that for $a_1 = a_2 = \frac{1}{2}$ the method is equal to the modified Euler method, and remember that this method is stable for real eigenvalues if $\Delta t \leq \frac{-2}{\lambda}$)

2. (a) The Taylor polynomials around 0 are given by:

$$f(0) = f(0) ,$$

$$f(-h) = f(0) - hf'(0) + \frac{h^2}{2}f''(0) - \frac{h^3}{6}f'''(\xi_1) ,$$

$$f(-2h) = f(0) - 2hf'(0) + 2h^2f''(0) - \frac{(2h)^3}{6}f'''(\xi_2) .$$

Here $\xi_1 \in (-h, 0)$, $\xi_2 \in (-2h, 0)$. We know that $Q(h) = \frac{\alpha_0}{h^2}f(0) + \frac{\alpha_{-1}}{h^2}f(-h) + \frac{\alpha_{-2}}{h^2}f(-2h)$, which should be equal to $f''(0) + \mathcal{O}(h)$. This leads to the following conditions:

$$\begin{array}{rcl} f(0): & \frac{\alpha_0}{h^2} & + & \frac{\alpha_{-1}}{h^2} & + & \frac{\alpha_{-2}}{h^2} & = & 0 \\ f'(0): & & -\frac{h\alpha_{-1}}{h^2} & - & \frac{2h\alpha_{-2}}{h^2} & = & 0 \\ f''(0): & & \frac{h^2}{2h^2}\alpha_{-1} & + & \frac{2h^2\alpha_{-2}}{h^2} & = & 1 \\ \end{array}$$

This can also be written as

$$\begin{array}{rcl} f(0): & \alpha_0 & + & \alpha_{-1} & + & \alpha_{-2} & = & 0 \\ f'(0): & & -\alpha_{-1} & - & 2\alpha_{-2} & = & 0 \\ f''(0): & & \frac{\alpha_{-1}}{2} & + & 2\alpha_{-2} & = & 1 \\ \end{array}$$

(b) The *truncation error* follows from the Taylor polynomials:

$$f''(0) - Q(h) = f''(0) - \frac{f(0) - 2f(-h) + f(-2h)}{h^2} = -\left(\frac{\frac{2h^3}{6}f'''(\xi_1) - \frac{8h^3}{6}f'''(\xi_2)}{h^2}\right)$$
$$= hf'''(\xi).$$

(c) Note that

$$f''(0) - Q(2h) = K2h \tag{21}$$

$$f''(0) - Q(h) = Kh$$
 (22)

Subtraction gives:

$$Q(h) - Q(2h) = K2h - Kh = Kh.$$
(23)

We choose $h = \frac{1}{4}$. Then $Q(2h) = Q(\frac{1}{2}) = \frac{0-2 \times 0.1250+1}{0.25} = 3$ and $Q(h) = Q(\frac{1}{4}) = \frac{0-2 \times 0.0156+0.1250}{(\frac{1}{4})^2} = 1.5008$. Combining (22) and (23) shows that

$$f''(0) - Q(\frac{1}{4}) = Q(\frac{1}{4}) - Q(\frac{1}{2}) = -1.4992$$

3. (a) First, we check the boundary conditions:

$$u(0) = 0 - \frac{1 - e^{0/\epsilon}}{1 - e^{1/\epsilon}} = \frac{1 - 1}{1 - e^{1/\epsilon}} = 0, \quad u(1) = 1 - \frac{1 - e^{1/\epsilon}}{1 - e^{1/\epsilon}} = 0.$$
(24)

Further, we have

$$u'(x) = 1 + \frac{e^{x/\epsilon}}{\epsilon(1 - e^{1/\epsilon})},$$
 (25)

$$u''(x) = \frac{e^{x/\epsilon}}{\epsilon^2 (1-e^{1/\epsilon})}.$$
 (26)

Hence, we immediately see

$$-\epsilon u''(x) + u'(x) = -\frac{\epsilon e^{x/\epsilon}}{\epsilon^2 (1 - e^{1/\epsilon})} + 1 + \frac{e^{x/\epsilon}}{\epsilon (1 - e^{1/\epsilon})} = 1.$$
 (27)

Hence, the solution $u(x) = 1 - \frac{1 - e^{x/\epsilon}}{1 - e^{1/\epsilon}}$ satisfies the differential equation and the Dirichlet boundary conditions, and therewith u(x) is the solution to the boundary value problem (uniqueness can be demonstrated in a straightforward way, but this was not asked for).

(b) The domain of computation, being (0, 1), is divided into subintervals with mesh points, we set $x_j = j\Delta x$, where we use *n* unknowns, such that $x_{n+1} = (n + 1)\Delta x = 1$. The discretization method for an interior node is given by

$$-\epsilon \frac{w_{j+1} - 2w_j + w_{j-1}}{(\Delta x)^2} + \frac{w_j - w_{j-1}}{\Delta x} = 1, \text{ for } j \in \{1, \dots, n\}.$$
 (28)

At the boundaries, we see for j = 1 and j = n, upon substituting $w_0 = 0$ and $w_{n+1} = 0$, respectively:

$$-\epsilon \frac{w_2 - 2w_1 + 0}{(\Delta x)^2} + \frac{w_1 - 0}{\Delta x} = 1,$$

$$-\epsilon \frac{0 - 2w_n + w_{n-1}}{(\Delta x)^2} + \frac{w_n - w_{n-1}}{\Delta x} = 1.$$
(29)

This can be rewritten more neatly as follows:

$$\epsilon \frac{-w_2 + 2w_1}{(\Delta x)^2} + \frac{w_1}{\Delta x} = 1,$$
(30)
$$\epsilon \frac{2w_n - w_{n-1}}{(\Delta x)^2} + \frac{w_n - w_{n-1}}{\Delta x} = 1.$$

(c) Next, we use $\Delta x = 1/4$ and thus n = 3. Then, from equations (28) and (30), one obtains the following 3×3 linear system of equations

$$\begin{bmatrix} 32\epsilon + 4 & -16\epsilon & 0\\ -16\epsilon - 4 & 32\epsilon + 4 & -16\epsilon\\ 0 & -16\epsilon - 4 & 16\epsilon + 4 \end{bmatrix} \begin{bmatrix} w_1\\ w_2\\ w_3 \end{bmatrix} = \begin{bmatrix} 1\\ 1\\ 1 \end{bmatrix}$$
(31)

(d) No, the upwind difference method for the convective terms is designed to not produce oscillatory solutions independent of the size of the diffusion coefficient ϵ and the velocity v which is equal to one in this case.